

Magnetic Reconnection in Slab Geometry



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2013 BOUT++ Workshop

Livermore, CA

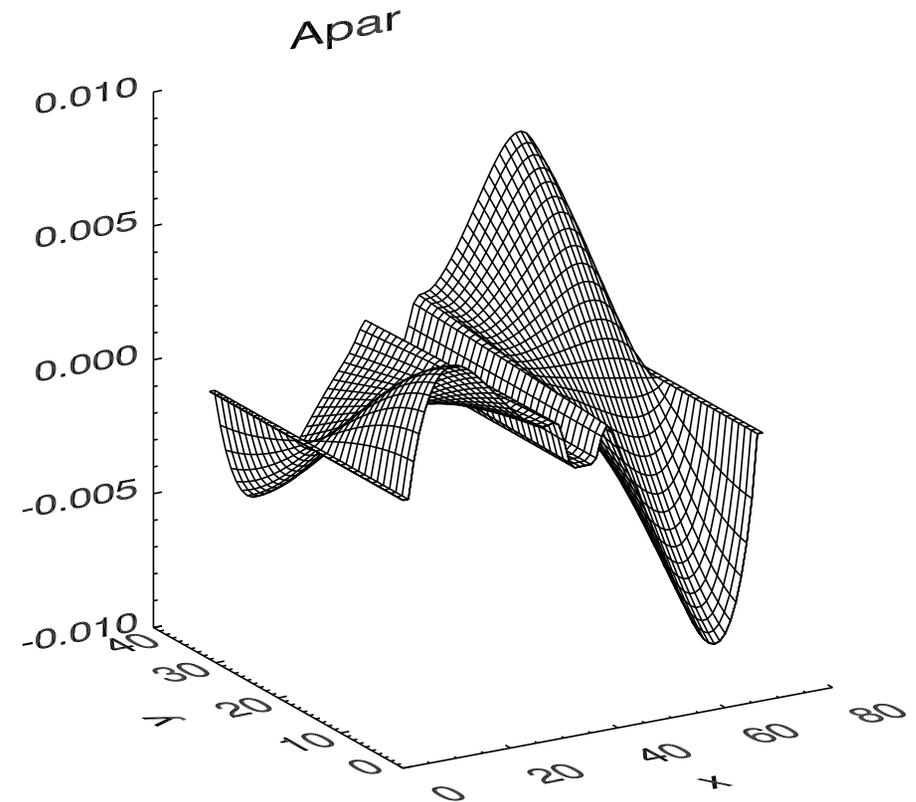
September 6th, 2013

**This work performed under the auspices of the U.S. Department of Energy by
Lawrence Livermore National Laboratory under Contract DE-AC52-07NA27344**

LLNL-PRES-643597

Outline

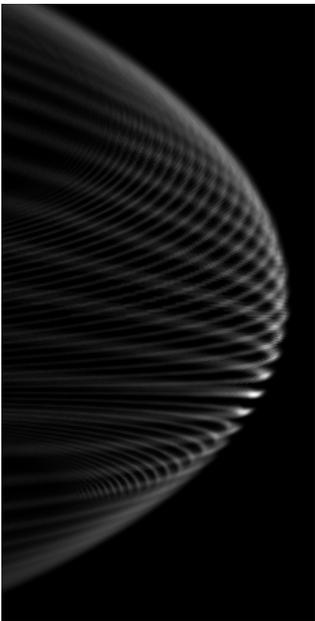
- Code overview
 - some recent improvements
 - slab grid generator
 - reconnect-2field
 - post-processing
- Benchmarking reconnection dynamics
 - Scaling with dissipation
 - Suppression due to rotation



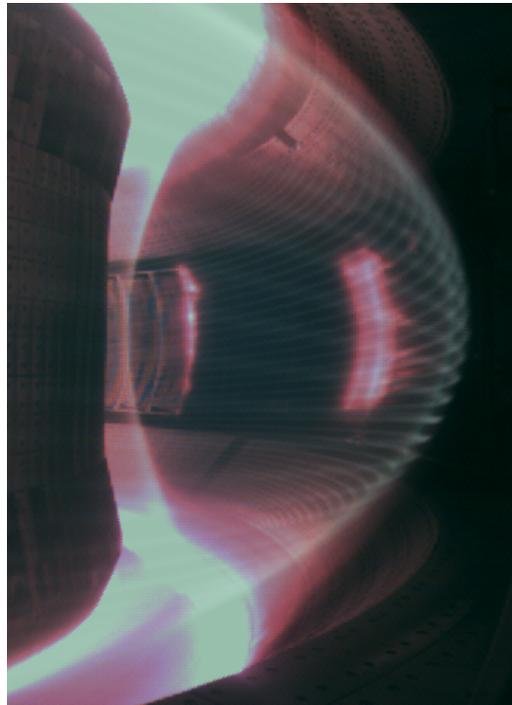
Some improvements to IDL tools since 2011

- I developed some tools to generate output to Bill Meyer's raytracing code
 - `write_polslice`, `resize_bgrid`

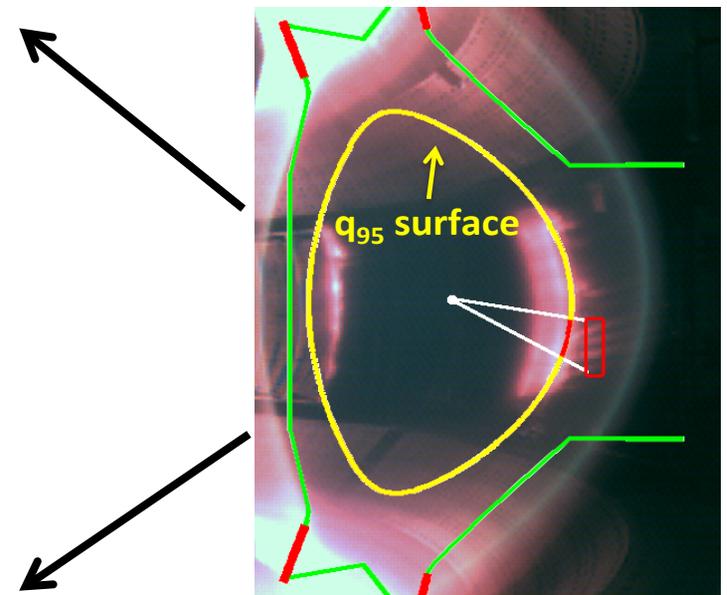
Simulation



Comparison



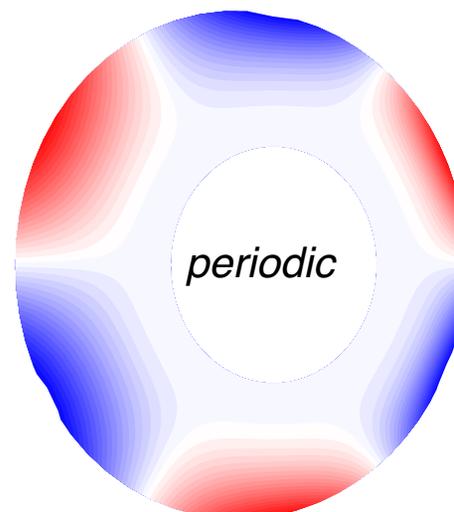
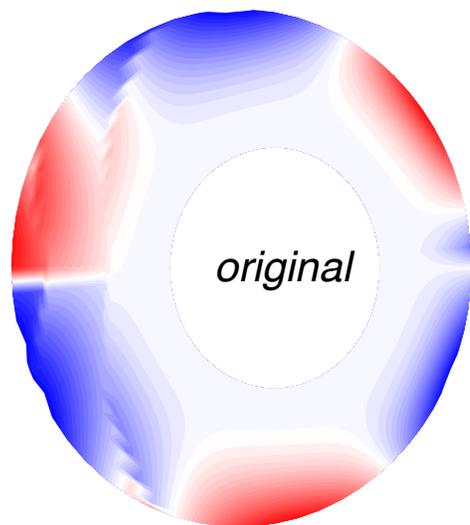
Measurement



T.Y. Xia, et al., 24th IAEA FEC, San Diego, Oct. 2012

Some improvements to IDL processing

- IDL processing improvements
 - Added some useful numerical utilities to `tools/numlib`
 - Clean up coding by using `CASE` to test dimension of input variable
 - Fixed error in `zshift` and `plotpolslice` for periodic domains
 - Please contribute!



Some improvements to core code

- **Core BOUT++ improvements (B. Dudson)**
 - Physics-based **preconditioning** for shear Alfvén wave greatly improves convergence for large time steps
 - **Must carefully set error tolerances to ensure that output is not **garbage****
 - ATOL and RTOL
 - New Option **SymmetricGlobalX = true** → initial profiles are symmetric in X
 - **otherwise there is an asymmetry in GLOBALX function!**
 - **Laplace inversion** with nonconstant BCs now gives correct result when **NXPE>1**
 - **allows simulations with more processors**

“Taylor’s problem” for resistive reconnection proceeds in a number of different temporal phases*

- **2 Field model**

$$\partial_t A_{\parallel} = -\nabla_{\parallel} \phi - \eta J_{\parallel}$$

$$J_{\parallel} = -\nabla_{\perp}^2 A_{\parallel}$$

$$\partial_t U + \nabla \cdot V_E U = V_A^2 \nabla \cdot J_{\parallel} + \mu \nabla_{\perp}^2 U$$

$$U = mn \nabla_{\perp}^2 \varphi / B$$

- **Spatial scales***

$$s = a q' / q = a B'_y / B_y = q R / R_s \sim 3.3$$

$$k_y \delta_{Rec} = (\tau_H / \tau_R)^{1/3} = S_R^{-1/3}$$

$$R_s = B_z / B'_y = k_y / k'_{\parallel} = q R / s \sim 4.5 m$$

$$k_y \delta_{RI} = (\omega \tau_H^2 / \tau_R)^{1/4} = S_R^{-1/4} S_{\omega}^{-1/4}$$

- **Intrinsic & derived time scales**

$$\tau_A = R_s / V_A = 0.64 \mu s$$

$$\tau_{Rec} = \tau_R^{1/3} \tau_A^{2/3} \propto \eta^{-1/3} k_y^{-2/3}$$

$$\tau_R = a^2 / \eta = 4.6 ms$$

$$\tau_{RI} = 1.8 \tau_R^{3/5} \tau_H^{2/5} \propto \eta^{-3/5} k_y^{-2/5}$$

$$\tau_V = a^2 / \mu$$

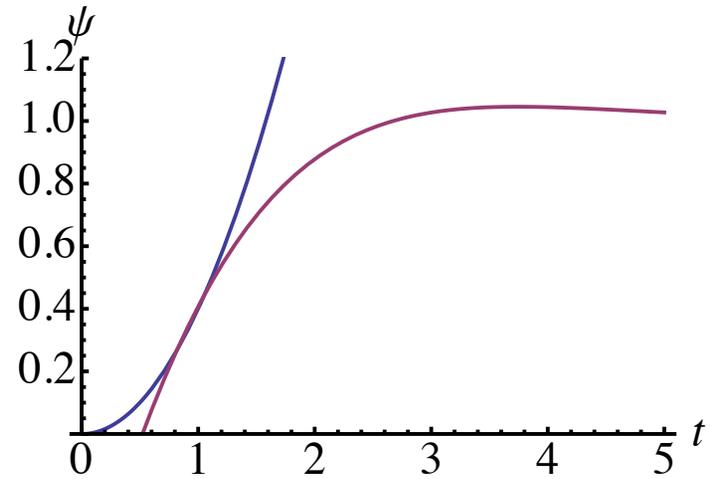
Linear “Taylor’s problem” proceeds in 2 distinct phases*

- Approximate linear analytic solution**

$$\psi \approx \begin{cases} At^2 / 2\tau_R\tau_A & t < \tau_{rec} \\ B\left(1 - \frac{8}{5}e^{-tc/\tau_{RI}} \cos(ts/\tau_{RI})\right) & t > \tau_{rec} \\ c + is = e^{i\pi/5} \end{cases}$$

$$\tau_{Rec} = \tau_R^{1/3} \tau_A^{2/3} \propto \eta^{-1/3} k_y^{-2/3}$$

$$\tau_{RI} = 1.8\tau_R^{3/5} \tau_H^{2/5} \propto \eta^{-3/5} k_y^{-2/5}$$

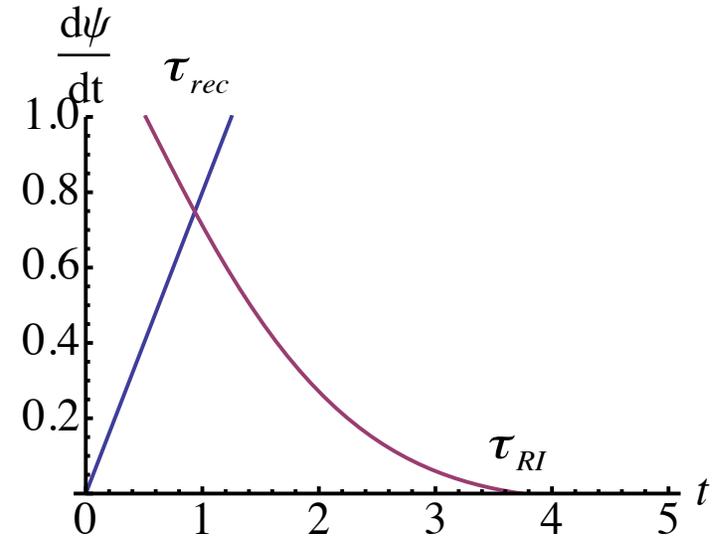


- Different scalings for

- time of peak rate = reconnection time τ_{Rec}

- time of peak flux = saturation time τ_{RI}

$$\max(\tau_A \dot{\psi}) \sim \tau_A \dot{\psi}(\tau_{Rec}) \propto \frac{\tau_{Rec}}{\tau_R} = \left(\frac{\tau_A}{\tau_R}\right)^{2/3} \propto \left(\frac{\eta}{k_y}\right)^{2/3}$$



*T.S. Hahm and R. M. Kulsrud, Phys. Fluids **28** (1985) 2412

A. Cole and R. Fitzpatrick, Phys. Plasmas **11 (2004) 3525

Slab geometry

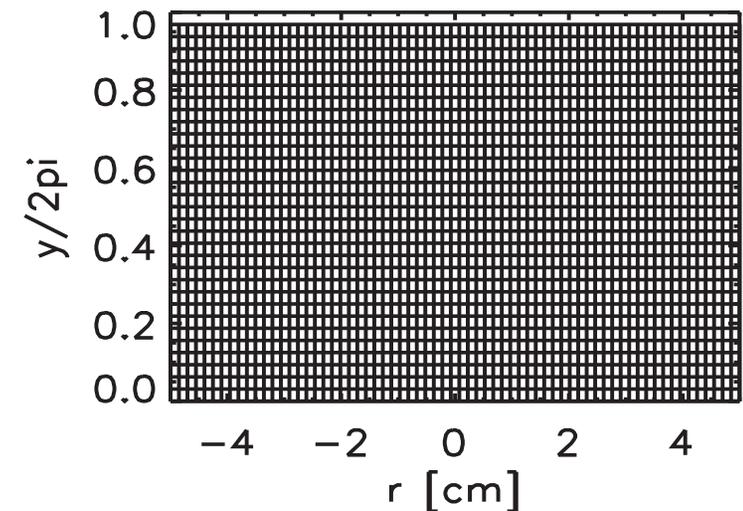
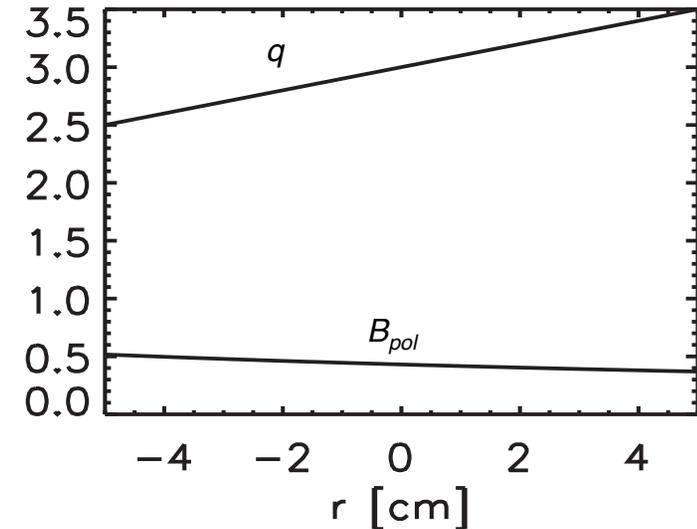
- Equilibrium

- $R = 5 \text{ m}$, $a = 1 \text{ m}$, $\varepsilon = 0.2$
- $B_{\text{tor}} = 6.46 \text{ T}$, $n_e = 10^{20} \text{ m}^{-3}$
- $B_{\text{pol}}/B_{\text{tor}} = \varepsilon/q = 6-8 \times 10^{-2}$
- $q = [2.5, 3.5]$
- $s = aq'/q = 10/3 \sim 3.3$, $L_s = qR/s = 4.5 \text{ m}$
- $\Delta r = 0.1 \text{ m}$, $\Delta\psi = 0.21 \text{ Vs}$
- $\rho = 1 \text{ mm}$, $\rho/\Delta r = 1/100$
- $\delta r = 0.1 \text{ m}/68 \sim 1.5 \text{ mm} \sim 1.5 \rho$

- BOUT mesh uses twist-shift periodic coordinates

- Given orthogonal coordinates: ψ , θ , ξ
- BOUT coordinates are $x = \psi/\delta\psi$, $y = \theta$, and

$$z = \xi - \int d\theta \frac{B \cdot \nabla \xi}{B \cdot \nabla \theta} = \xi - q\theta + \text{periodic}$$



Slab Grid Generator: \$BOUT_TOP/tools/slab/slab.pro

```
tools/slab> more slab.pro
; Slab geometry grid generator
;
; Optional keywords:
; =====
;
; output   = Set output file name
; thin     = Use thin radial box approximation
;           so Bpxy = constant, but gradient is non-zero
;
; ni       = Ion density in 10^20 m^-3
; Ti       = temperature in eV (Te = Ti)
;
; Rmaj     = Major radius [meters]
; rminor   = Minor radius [m]
; dr       = Radial width of box [m]
; r_wid    = Radial extent, normalised to gyro-radius r_wid = dr / rho_i
;
; q        = Safety factor q = r*Bt/(R*Bp) at middle of box
; dq       = Change in q. Will go from q-dq/2 to q+dq/2
;
; L_T      = Temperature length scale [m]
; eta_i    = Ratio of density to temp. length scales eta = L_n / L_T

PRO slab, output=output, thin=thin, $
      ni=ni, Ti=Ti, $
      Rmaj=Rmaj, rminor=rminor, dr=dr, $
      r_wid=r_wid, $
      q=q, dq=dq, $
      L_T=L_T, eta_i=eta_i, $
      nx=nx, ny=ny
```

Generate a Slab Grid

```
tools/slab> more idlsave.pro
; IDL Version 8.2.2 (linux x86_64 m64)
; Journal File for train213@hopper09
; Working directory: /global/u2/t/train213/BOUT_Workshop_2013/BOUT-2.0/to
b
; Date: Wed Sep  4 14:12:39 2013

nx=2^8+4 & print, nx
;    260
ny=64
file = 'slab_'+str(nx)+'x'+str(ny)+'.nc' & print, file
;slab_260x64.nc
slab, Rmaj=10,rmin=1,dr=.1, nx=nx, ny=ny, out=file
;Poloidal field varies from 0.184642 to 0.258498
;Writing grid to file slab_260x64.nc
;DONE
g=file_import(file)
help, g
tools/slab> █
```

2-field model: \$BOUT/examples/reconnect-2field/2field.cxx

```
examples/reconnect-2field> cd $BOUT_TOP/examples/reconnect-2field
examples/reconnect-2field> make clean
examples/reconnect-2field> make
  Compiling 2field.cxx
  Linking 2field
examples/reconnect-2field> ls
2field      2field-new.cxx  2field-precon.cxx  data      makefile-new
2field.cxx  2field.o       BOUT-new.inp      makefile  slab_68x32.nc
examples/reconnect-2field> more 2field.cxx
/*****
 * 2 field (Apar, vorticity) model for benchmarking
 * simple slab reconnection model
 *****/

#include <bout.hxx>
#include <boutmain.hxx>

#include <invert_laplace.hxx>
#include <invert_parderiv.hxx>
#include <initialprofiles.hxx>
#include <bout/constants.hxx>

#include <math.h>
#include <stdio.h>
#include <stdlib.h>

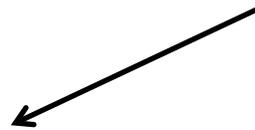
// 2D initial profiles
Field2D Jpar0, Te0, Ni0;

// 3D evolving fields
Field3D Upar, Apar;

// Derived 3D variables
Field3D Phi, Jpar;

// External fields
Field3D Apar_ext, Jpar_ext, Phi0_ext, Upar0_ext;
```

**External Fields
Stay Constant
& Don't Evolve**



Apply fixed external field A_{ext}

- Chosen form is resonant for $q=3$ in center of domain

$$A_{||} = A_{edge} (1 - 4x(1 - x)) \sin(3\theta - \xi)$$

- Radial profile chosen to vanish at the center to avoid reconnection at $t=0$

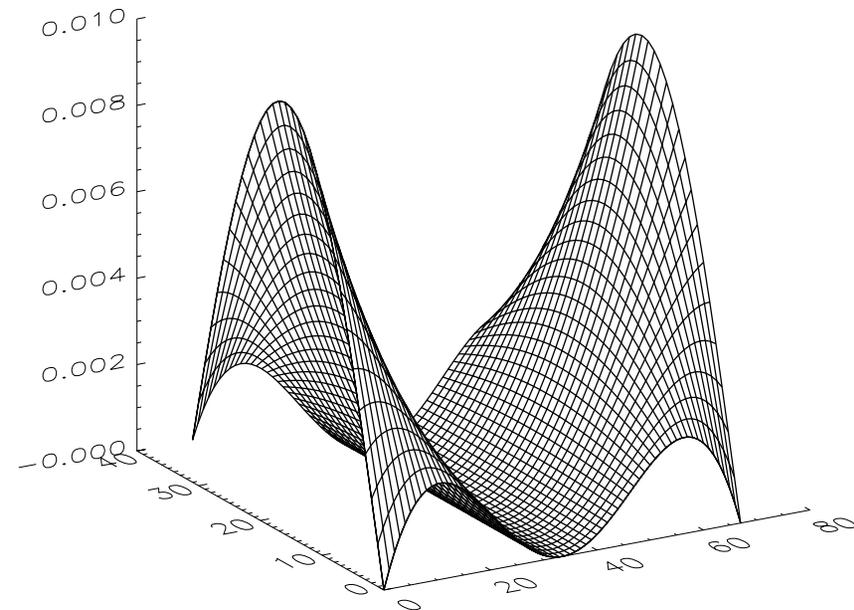
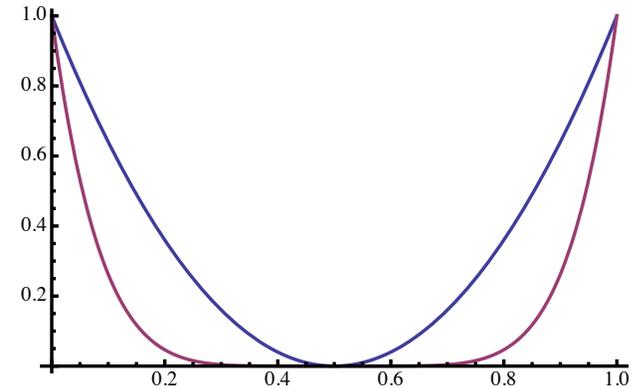
- 2field code treats variables with ordering:

$$A_{||} = \delta(A_{ext} + A_1)$$

$$J_{||} = -\nabla_{\perp}^2 A_{||}$$

$$\varphi = \varphi_{0,ext} + \delta\varphi_1$$

$$U_1 = \nabla_{\perp}^2 \varphi$$



Run a 2field example

```
examples/reconnect-2field> qsub -I -V -q interactive -l mppwidth=2 -l advres=bout,10
qsub: waiting for job 6436873.hopque01 to start
qsub: job 6436873.hopque01 ready

ModuleCmd_Switch.c(172):ERROR:152: Module 'PrgEnv-pgi' is currently not loaded
t/train213> cd $BOUT_TOP/examples/reconnect-2field/
examples/reconnect-2field> aprun -n 2 ./2field

BOUT++ version 1.10
Revision: 546769c683863172c675d673f10c106ff9f5a2dc
Code compiled on Sep  4 2013 at 10:15:12

B.Dudson (University of York), M.Umansky (LLNL) 2007
Based on BOUT by Xueqiao Xu, 1999

Processor number: 0 of 2

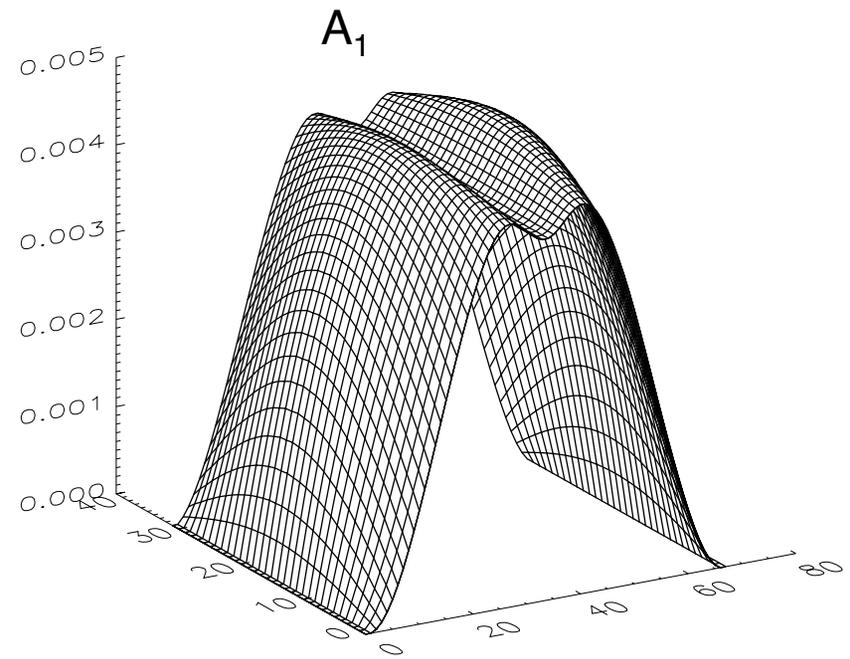
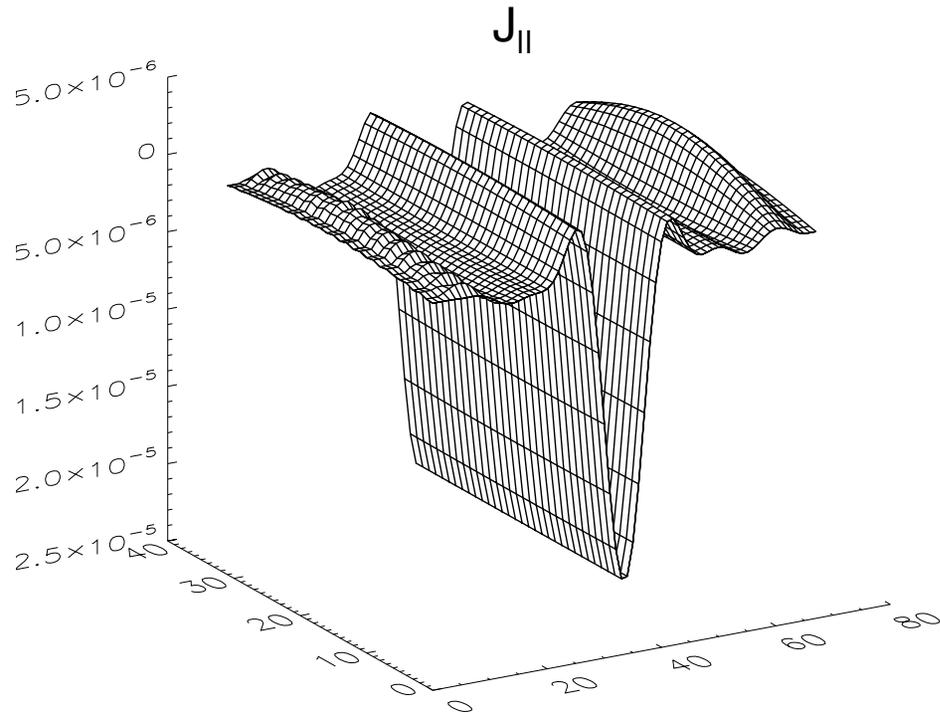
pid: 19920

Compile-time options:
  Checking enabled, level 2
  Signal handling enabled
  PDB support disabled
  netCDF support enabled
  Parallel NetCDF support disabled
  OpenMP parallelisation disabled
Reading options file data/BOUT.inp
  Option /append = false (default)
  Option /dump_format = nc (data/BOUT.inp)
  Option /staggergrids = false (data/BOUT.inp)
```

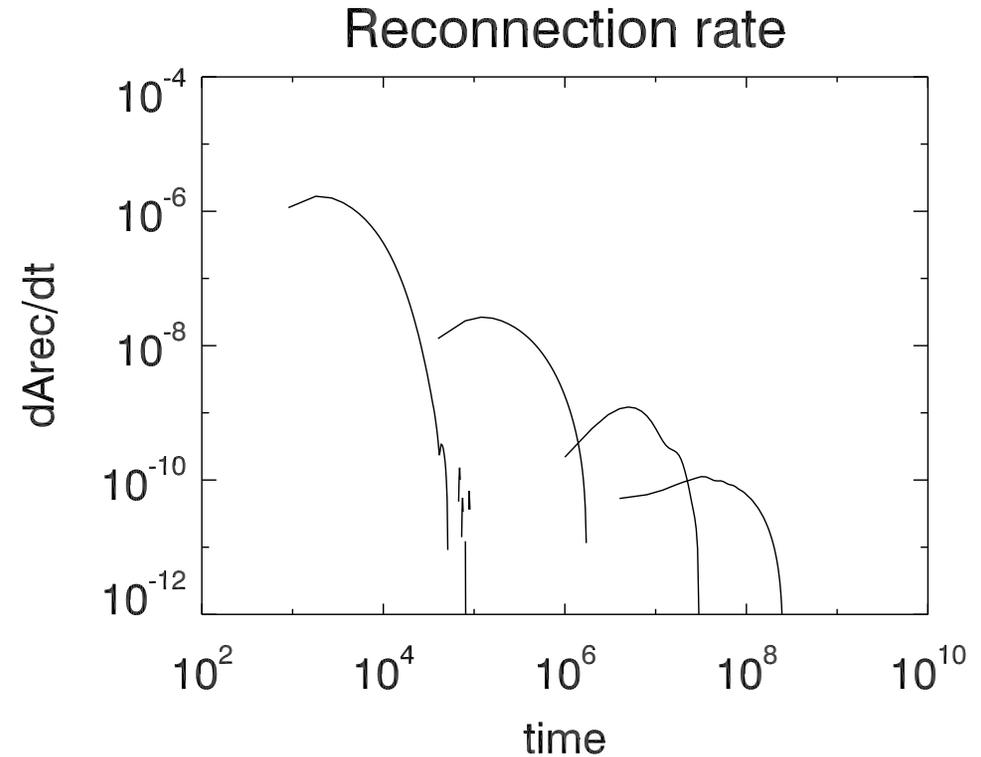
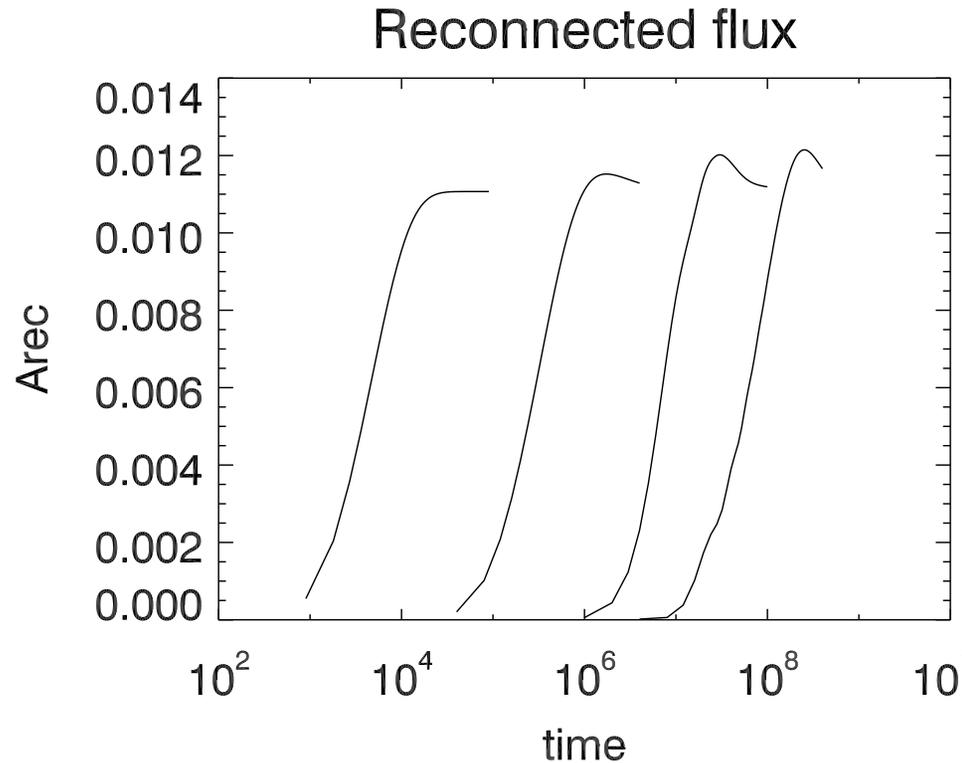
Examine data

```
aext=collect(var='Apar_ext',path='data')
surface, reform(aext[:,*,4,:])
apar=collect(var='Apar',path='data')
showdata, reform(apar[:,*,4,:])
;chars=      2
```

```
examples/reconnect-2field> █
```

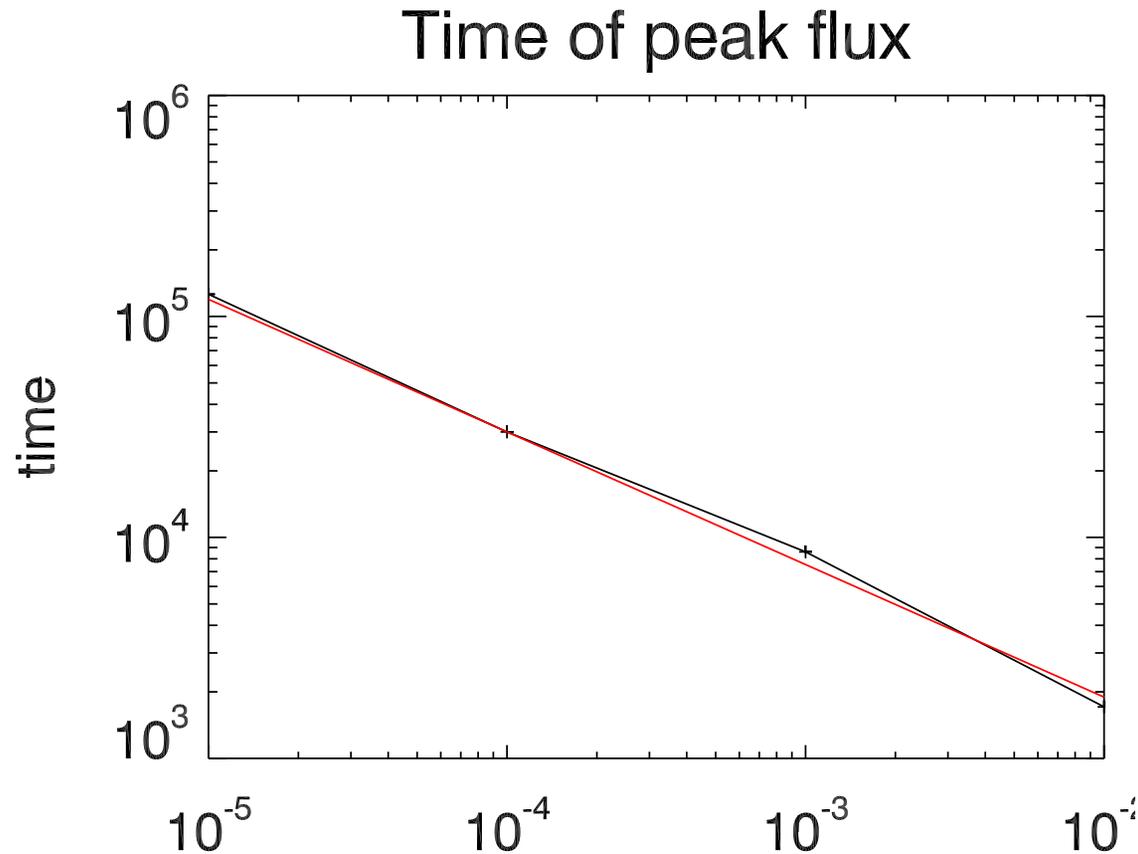


Evolution of reconnected flux for resistive model



- Reconnected flux generally peaks & relaxes as expected
- Initial growth is quadratic in time: t^2
- Resistivity range $\eta/D_B = 10^{-6}-10^{-4}$

Intrinsic time-scales demonstrate theoretical resistivity scaling



- *Time of peak flux displays resistive-inertial reconnection time: $\tau \sim \eta^{-3/5}$*

Linear reconnection theory: visco-resistive MHD

reviewed in Fitzpatrick PoP (1998)

- **2-field plasma model** (at low beta)

$$\partial_t A + \nabla_{\parallel} \phi = \eta \nabla_{\perp}^2 A$$

$$(\partial_t + v_E \cdot \nabla) \nabla_{\perp}^2 \phi = -V_A^2 \nabla_{\parallel} \nabla_{\perp}^2 A + \mu_i \nabla_{\perp}^4 \phi$$

$$\mu_s \rightarrow \mu_s / m_s n_s$$

$$\eta \rightarrow \eta_{\parallel} / \mu_0$$

- **2 reconnecting cases:** constant ψ
- **2 ideal cases:** not constant $\psi(r_s)=0$

Regime	Δ	δ	$\gamma \tau_A$
RI: inertia + resistivity	$2.1 S_R^{3/4} \Gamma^{5/4}$	$(\Gamma / S_R)^{1/4}$	$0.64 S_R^{-3/5}$
VR: viscosity + resistivity	$2.1 S_R^{5/6} S_V^{-1/6} \Gamma$	$S_R^{-1/6} S_V^{-1/6}$	$0.48 S_R^{-5/6} S_V^{1/6}$
VI: viscosity + inductance	$4.6 S_V^{1/4} \Gamma^{-1/4}$	$(\Gamma / S_V)^{1/4}$	$450 S_V$
I: inertia + inductance	$\pi \Gamma^{-1}$	$1 / \Gamma$	$1 / \pi$

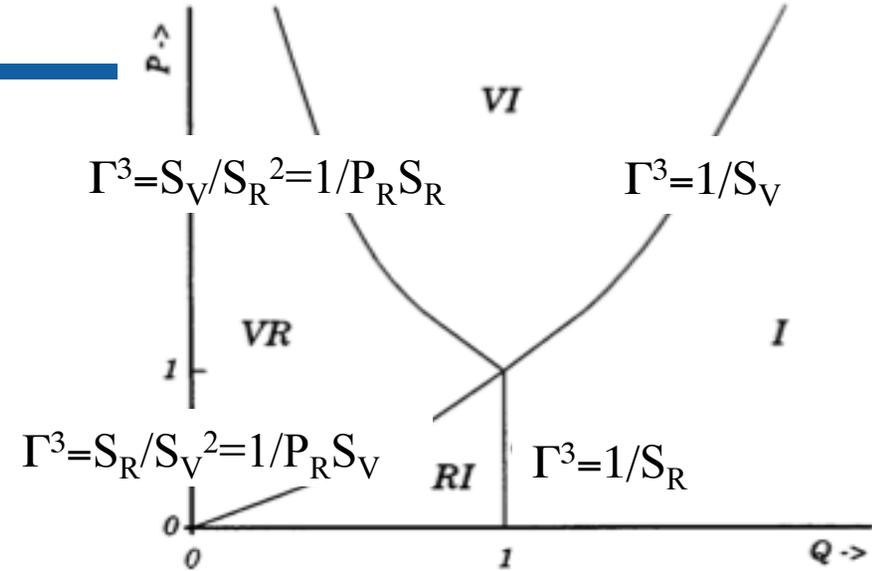


FIG. 1. A schematic diagram showing the extent of the four linear response regimes in normalized viscosity, P , versus normalized "slip frequency," Q , space. The four regimes are the visco-resistive regime (VR), the resistive-inertial regime (RI), the visco-inertial regime (VI), and the inertial regime (I).

- **Dimensionless #'s**

$$S_R = \tau_R / \tau_A = Lv_A / \eta$$

$$S_V = \tau_V / \tau_A = Lv_a / \mu_i$$

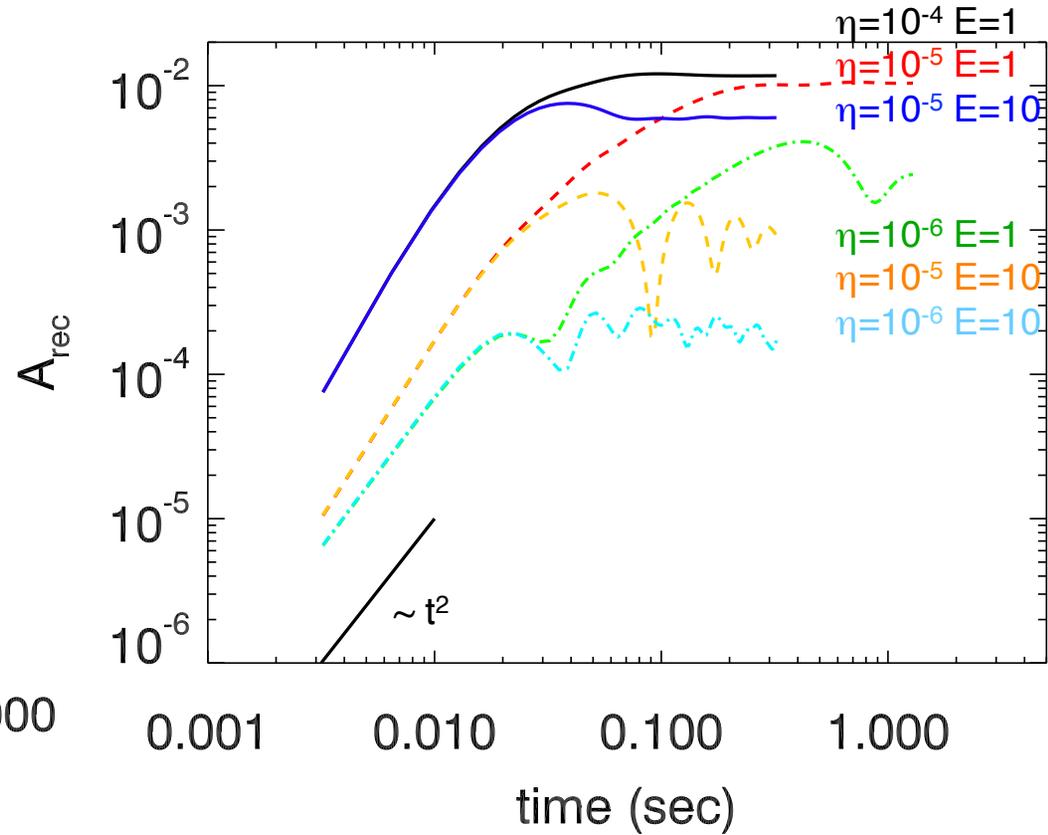
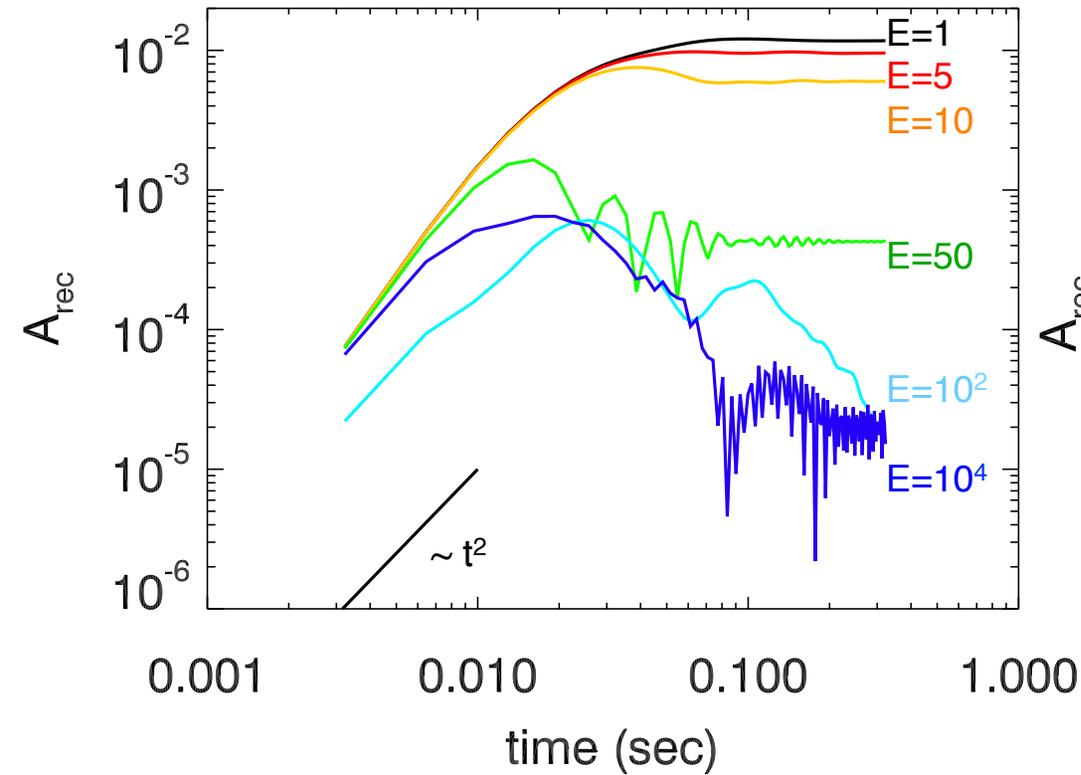
$$P_R = \tau_R / \tau_V = \mu_i / \eta$$

$$\Gamma = -i\omega\tau_A \quad Q_R = \omega\tau_A S_R^{1/3}$$

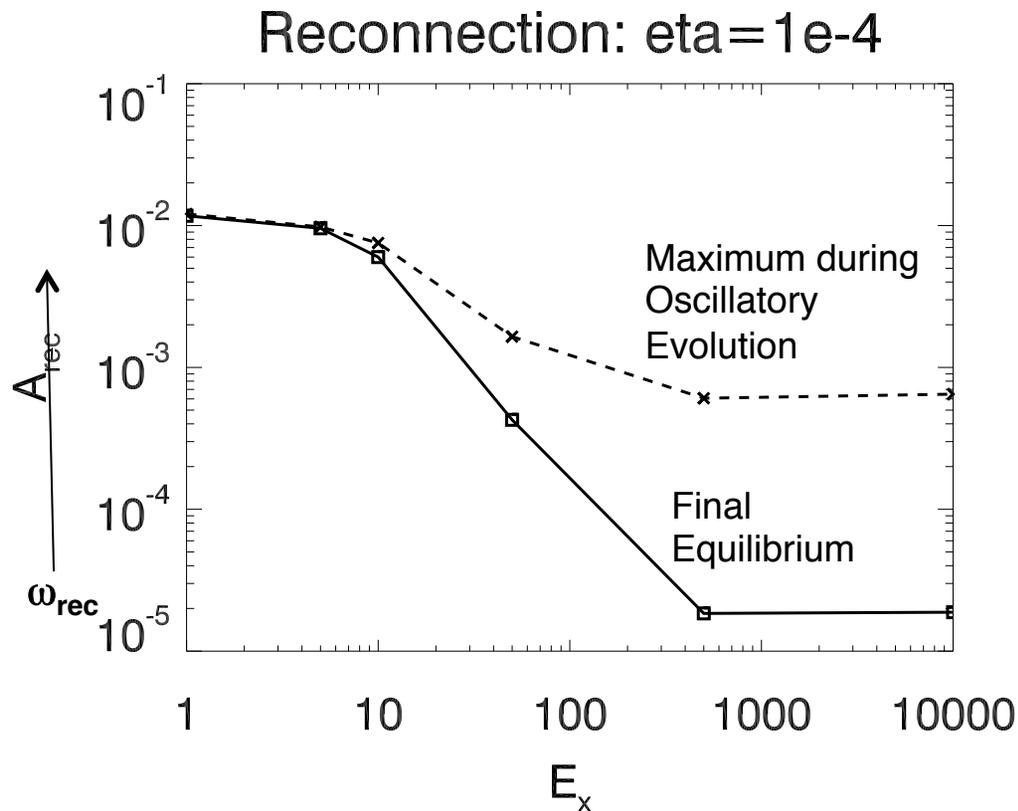


Perpendicular rotation suppresses reconnection: $\Phi_0 = -E_0 x$

$\eta = 10^{-4} D_B, S_R = 7 \times 10^8$

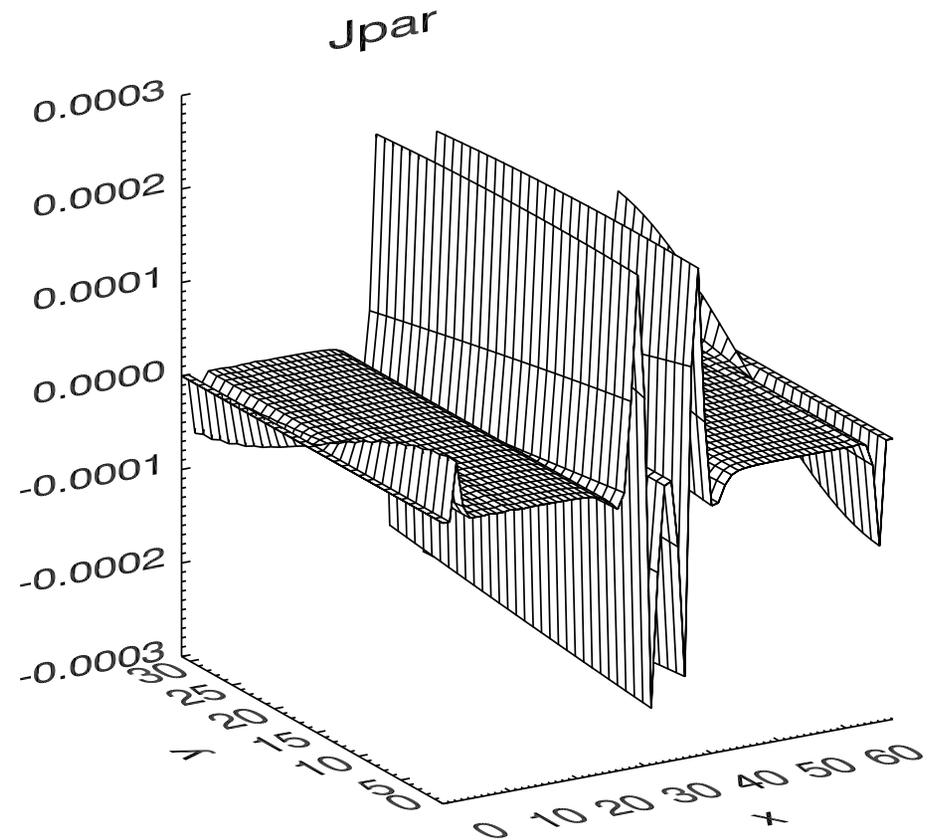
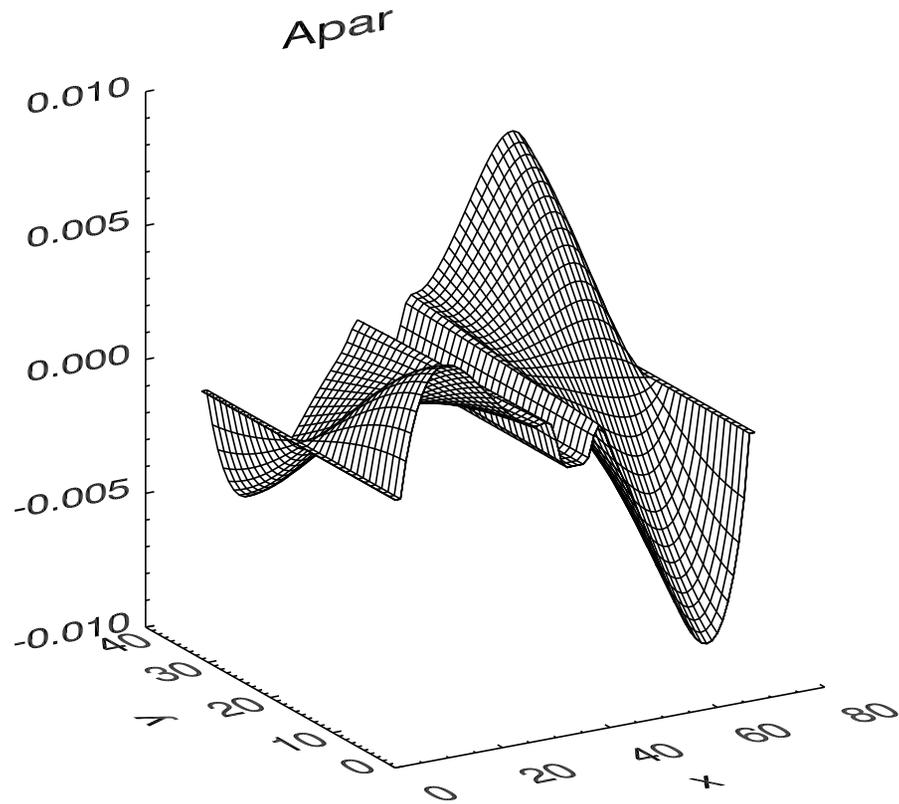


Suppression vs. Rotation: $\eta = 10^{-4} D_B$, $S_R = 7 \times 10^8$



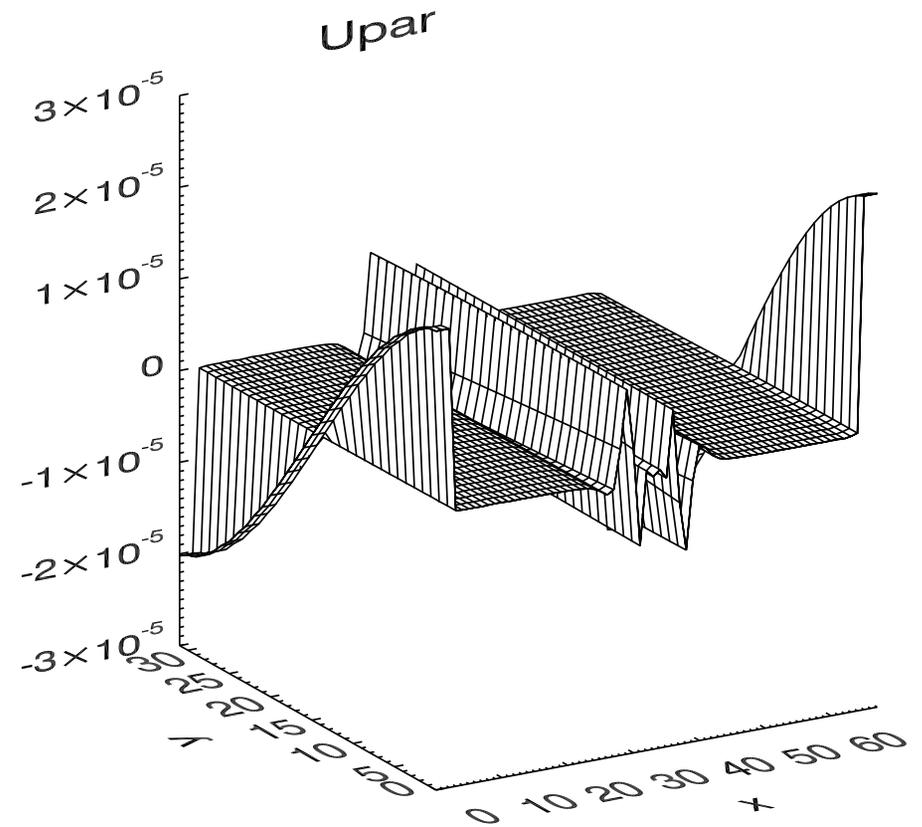
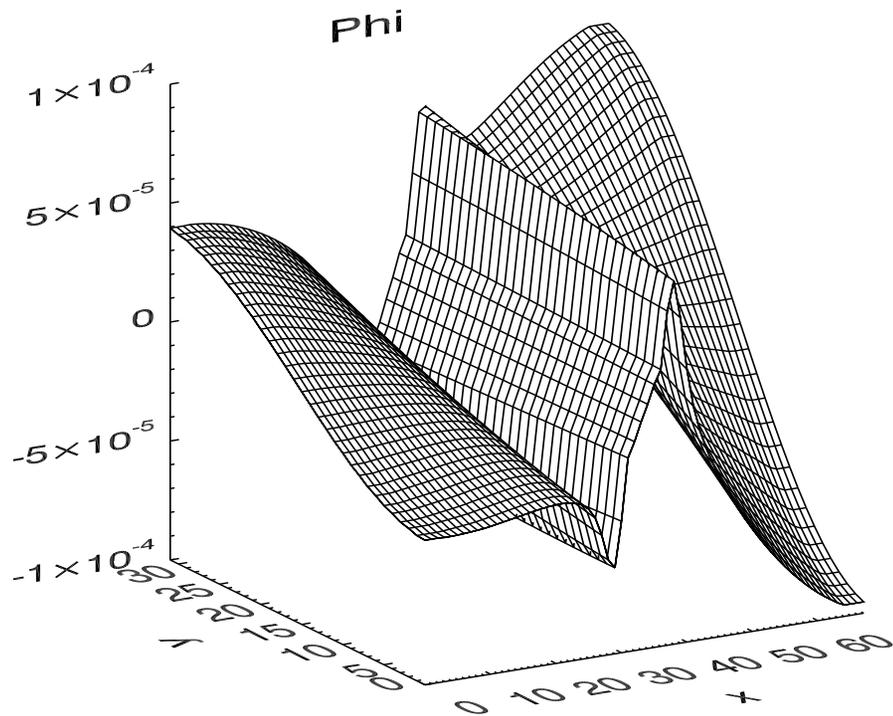
- Reconnection rotation rate measured to be $\omega_{rec} \sim 10 \omega_{E0} \sim 25$ krad/s
 - close to expected value $\omega_{rec} \sim S_R^{1/3} \omega_A \sim 23$ krad/s
- Phase difficult to measure accurately due to oscillations

Alfven resonance structure can be clearly seen at low resistivity & high rotation [$\phi=5, \eta=1e-7$]



- Alfvén resonance spacing $x_A \sim \omega_E / \omega_A k_y \sim 5.7 \text{ mm} \sim 4 \text{ grid points}$

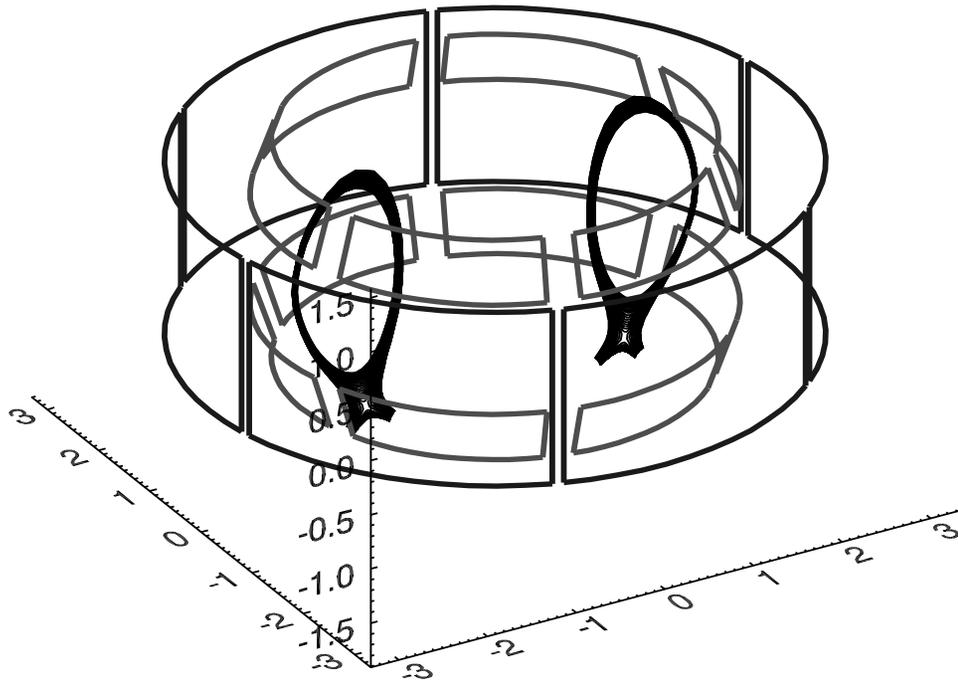
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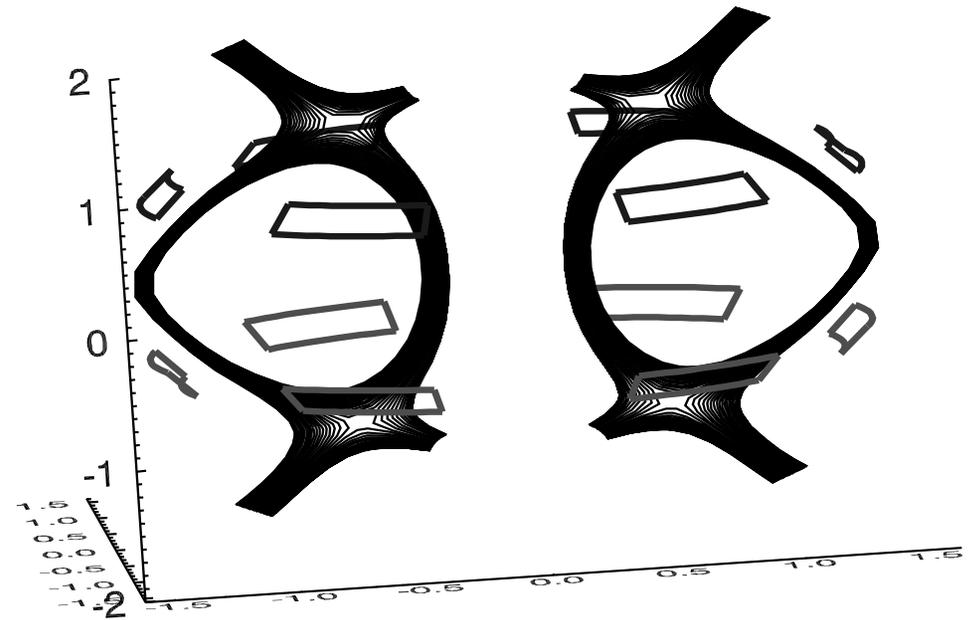
- Alfvén resonance spacing $x_A \sim \omega_E / \omega_A k_y \sim 5.7 \text{ mm} \sim 4 \text{ grid points}$

Moving toward realistic geometry: Implementation of 3D coil models improved

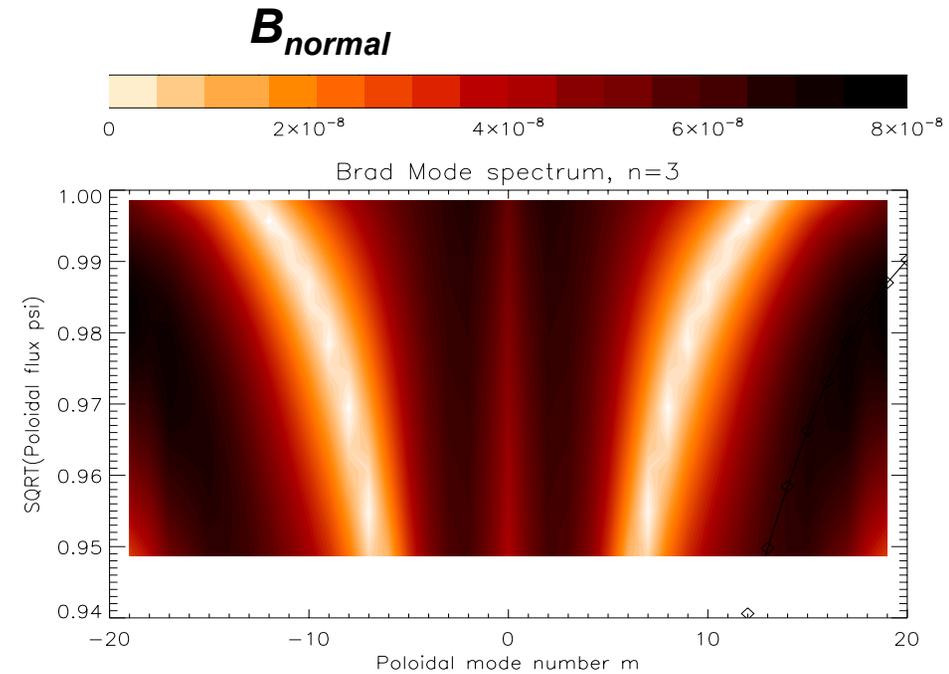
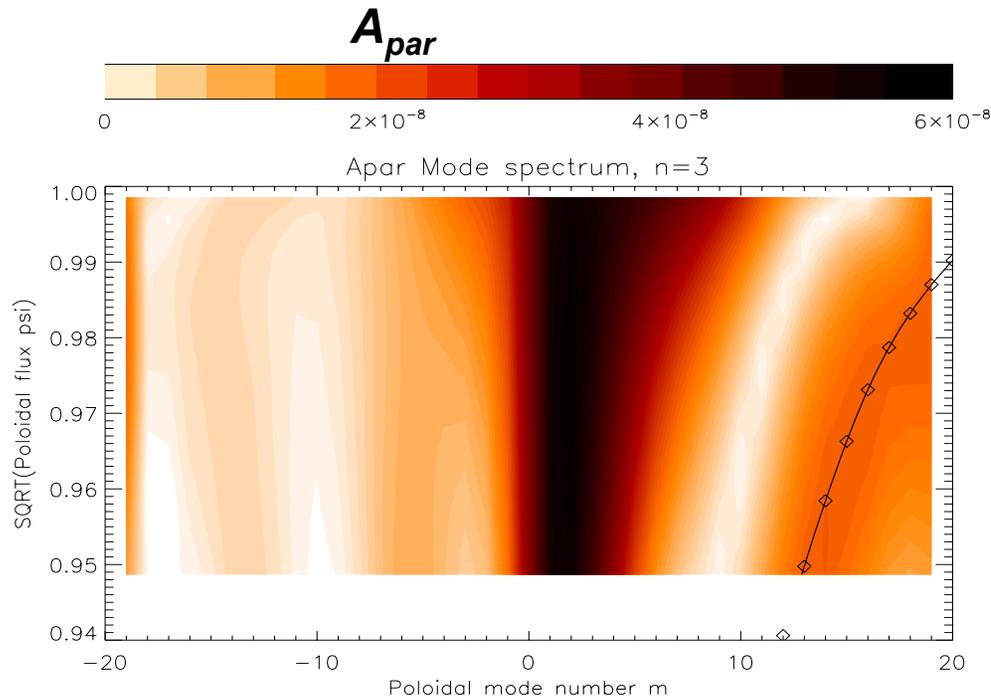
- DIII-D: I-coils & C-coils



- MAST: control coils



Magnetic field spectrum for MAST coil geometry



- Computed by representing coil fields as a superposition of fields due to individual line segments¹

¹Hanson & Hirshman, Phys. Plasmas 9, 4410 (2002)

Conclusions

- **Improvements made to allow robust treatment of BCs and coils**
 - **Realistic coil geometry implemented & will be benchmarked vs. other codes**

- **Linear & nonlinear reconnection benchmarks have been performed**
 - **Slab models of 2-field reconnection are being benchmarked vs. theory**
 - **Studies of hyper-resistive reconnection physics show different scaling laws**

- **Also: Ideal MHD boundary conditions have been implemented**
 - **Smoother relaxation to equilibrium, less generation of standing wave with large edge currents**

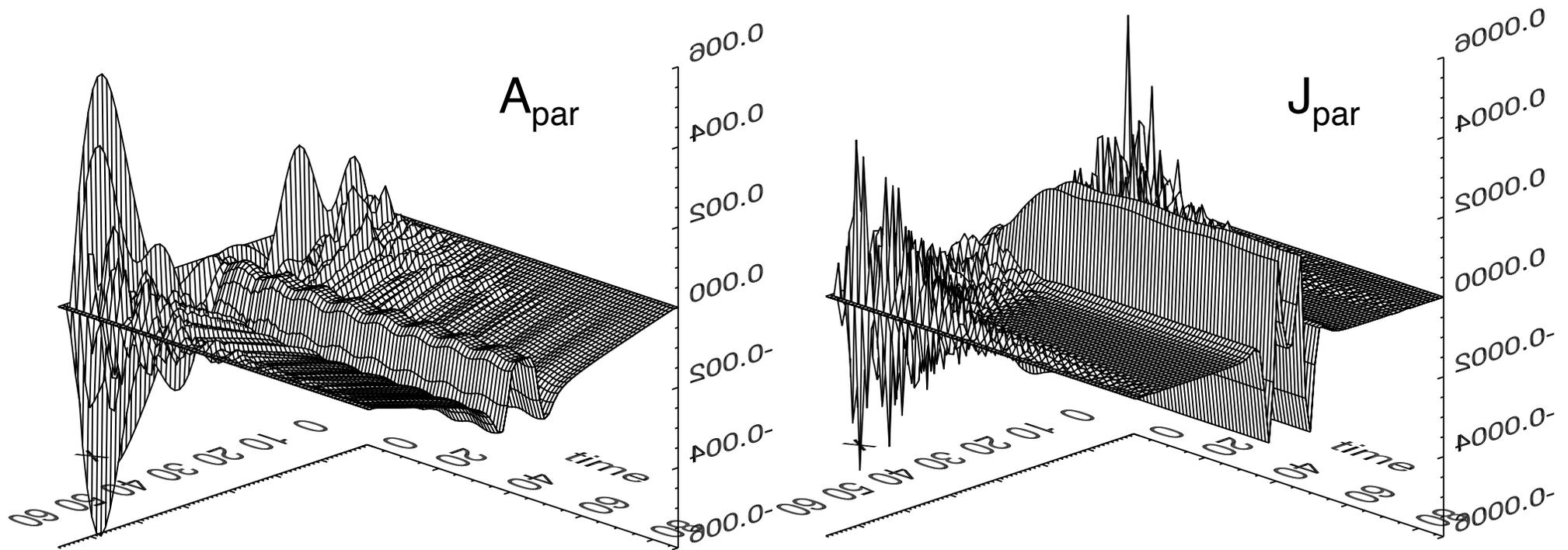
Exercises

- Make slab grid
 - `cd $BOUT_TOP/tools/slab`
 - `idl`
 - `slab, Rmaj=5, rmin=1, dr=0.1, nx=68, ny=32, output='myslab_68x32.nc'`
 - `exit`
- Run an example
 - `cd $BOUT_TOP/examples/reconnect-2field`
 - `git pull`
 - `make`
 - `>> edit data/BOUT.inp ; set parameters of interest`
 - `cp $BOUT_TOP/tools/slab/myslab_68x32.nc .`
 - `qsub -IV -q interactive -l mppwidth=16 -l advres=bout.10`
 - `cd $PBS_O_WORKDIR`
 - `aprun -n 16 ./2field -d data`
 - `idl ; examine data`
 - `apar_ext=collect(var='Apar_ext',path='data')`

Exercises

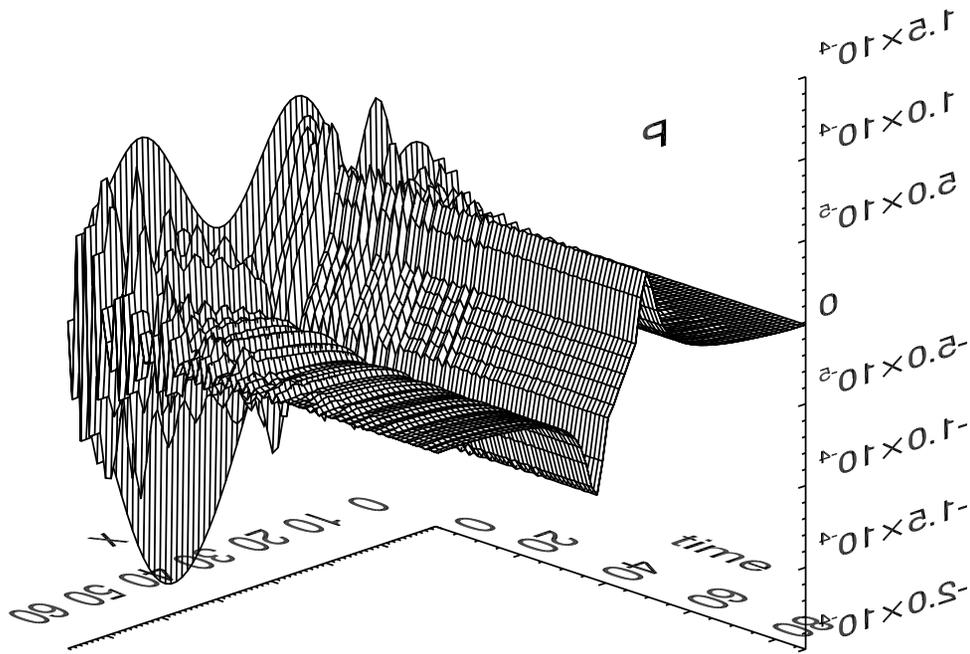
- Examine data
 - `idl ; examine data`
 - `apar_ext=collect(var='Apar_ext',path='data')`
 - `apar=collect(var='Apar',path='data')`
 - `jpar=collect(var='Jpar',path='data')`
 - `surface, reform(apar_ext[*,*],16)`
 - `showdata, reform(apar[*,*],16,*)`
- Run an example with preconditioning on/off
- Investigate behavior with changes to RTOL and ATOL

Alfven resonance structure can be clearly seen at low resistivity & high rotation [$\phi=5, \eta=1e-7$]



Alfven resonance structure can be clearly seen at low resistivity & high rotation [$\phi=5, \eta=1e-7$]

Phi



U_{par}

